

L^2 is a Hermitian operator

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A proof that square of the angular momentum vector is a Hermitian operator.

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We need to show that

$$\int \psi_1^* (L^2 \psi_2) d\Omega = \left[\int \psi_2^* (L^2 \psi_1) d\Omega \right]^*, \quad (1)$$

where

$$L^2 = - \left(\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \sin \theta \frac{\partial}{\partial \theta} + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right). \quad (2)$$

The second part of L^2 operator is easier to handle. The relevant part of the integral is the ϕ integral, which can be computed as follows

$$\begin{aligned} \int_0^{2\pi} \psi_1^* \left[\frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \psi_2 \right] d\phi &= \int_0^{2\pi} \frac{\partial}{\partial \phi} \left[\psi_1^* \frac{1}{\sin^2 \theta} \frac{\partial}{\partial \phi} \psi_2 \right] d\phi - \int_0^{2\pi} \left(\frac{\partial}{\partial \phi} \psi_1 \right)^* \left(\frac{1}{\sin^2 \theta} \frac{\partial}{\partial \phi} \psi_2 \right) d\phi \\ &= - \int_0^{2\pi} \frac{\partial}{\partial \phi} \left[\left(\frac{\partial}{\partial \phi} \psi_1 \right)^* \frac{1}{\sin^2 \theta} \psi_2 \right] d\phi + \int_0^{2\pi} \left(\frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \psi_1 \right)^* \psi_2 d\phi \\ &= \left[\int_0^{2\pi} \left(\frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \psi_1 \right) \psi_2^* d\phi \right]^*, \end{aligned} \quad (3)$$

where we dropped the first terms in the first two lines as they are the difference of the integrand at $\phi = 2\pi$ and $\phi = 0$, and that is zero as ϕ coordinate is 2π periodic.

The first part of L^2 operator seems to be harder because when we integrate by parts $\frac{\partial}{\partial \theta}$ will act on $\sin \theta$, which will complicate the problem. However, we can avoid it by a change of variable $u = \cos \theta$. The relevant part of the integral is the $d\cos \theta$ integral, and with the above transformation it becomes,

$$\begin{aligned}
\int_0^\pi \psi_1^* \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \psi_2 \right) d \cos \theta &= \int_{-1}^1 \psi_1^* \frac{\partial}{\partial u} \left((1-u^2) \frac{\partial}{\partial u} \psi_2 \right) du \\
&= - \int_{-1}^1 \frac{\partial}{\partial u} \psi_1^* \left((1-u^2) \frac{\partial}{\partial u} \psi_2 \right) du \\
&= - \int_{-1}^1 \left((1-u^2) \frac{\partial}{\partial u} \psi_1^* \right) \frac{\partial}{\partial u} \psi_2 du \\
&= \int_{-1}^1 \frac{\partial}{\partial u} \left((1-u^2) \frac{\partial}{\partial u} \psi_1^* \right) \psi_2 du \\
&= \left[\int_{-1}^1 \psi_2^* \frac{\partial}{\partial u} \left((1-u^2) \frac{\partial}{\partial u} \psi_1 \right) du \right]^* \\
&= \left[\int_0^\pi \psi_2^* \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \psi_1 \right) d \cos \theta \right]^*, \quad (4)
\end{aligned}$$

where we dropped again some surface terms as $u^2 - 1 = 0$ at $u = \pm 1$. (If you prefer $\sin \theta d\theta$ integral instead of $d \cos \theta$ integral, you will not need to change the variable.) This completes the proof that L^2 is Hermitian.

And that's it!